



Sampling algorithms in numerical linear algebra and their applications

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SLA-2014

Kalamata, Sept. 12, 2014

Introduction

- ‘Random Sampling’ or ‘probabilistic methods’: use of random data to solve a given problem.
- Eigenvalues, eigenvalue counts, traces, ...
- Many well-known algorithms use a form of random sampling: The Lanczos algorithm
- Recent work : probabilistic methods - See [Halko, Martinson, Tropp, 2010]
- Huge interest spurred by ‘big data’
- In this talk: A few specific applications of random sampling in numerical linear algebra

Introduction: A few examples

Problem 1: Compute $\text{Tr}[\text{inv}[A]]$ the trace of the inverse.

➤ Arises in cross validation :

$$\frac{\|(I - A(\theta))g\|_2}{\text{Tr}(I - A(\theta))} \quad \text{with} \quad A(\theta) \equiv I - D(D^T D + \theta L L^T)^{-1} D^T,$$

D == blurring operator and L is the regularization operator

➤ In [Huntchinson '90] $\text{Tr}[\text{Inv}[A]]$ is stochastically estimated

➤ Motivation for the work [Golub & Meurant, “Matrices, Moments, and Quadrature”, 1993, Book with same title in 2009]

Problem 2: Compute $\text{Tr} [f (A)]$, f a certain function

Arises in many applications in Physics. Example:

➤ Stochastic estimations of $\text{Tr} (f(A))$ extensively used by quantum chemists to estimate Density of States, see

[Ref: H. Röder, R. N. Silver, D. A. Drabold, J. J. Dong, Phys. Rev. B. 55, 15382 (1997)]

➤ Will be covered in detail later in this talk.

Problem 3: Compute $\text{diag}[\text{inv}(A)]$ the diagonal of the inverse

- Harder than just getting the trace
- Arises in Dynamic Mean Field Theory [DMFT, motivation for our work on this topic].
- Related approach: Non Equilibrium Green's Function (NEGF) approach used to model nanoscale transistors.
- In **uncertainty quantification**, the diagonal of the inverse of a covariance matrix is needed [Bekas, Curioni, Fedulova '09]

Problem 4: Compute $\text{diag}[f(A)]$; f = a certain function.

- Arises in any density matrix approach in quantum modeling - for example Density Functional Theory.
- Here, f = Fermi-Dirac operator:

$$f(\epsilon) = \frac{1}{1 + \exp\left(\frac{\epsilon - \mu}{k_B T}\right)}$$

Note: when $T \rightarrow 0$ then $f \rightarrow$ a step function.

Note: if f is approximated by a rational function then $\text{diag}[f(A)] \approx$ a linear combination of terms like $\text{diag}[(A - \sigma_i I)^{-1}]$

- **Linear-Scaling methods** based on approximating $f(H)$ and $\text{Diag}(f(H))$ – avoid ‘diagonalization’ of H

- Rich literature on 'linear scaling' or 'order n' methods
- The review paper [Benzi, Boito, Razouk, "Decay properties of Spectral Projectors with applications to electronic structure", SIAM review, 2013] provides theoretical foundations
- Several references on approximating $\text{Diag}(f(H))$ for this purpose – See e.g., work by L. Lin, C. Yang, E. E [Code: SellInv]
- Also: analysis of **network graphs** [yesterday's talk by Lothar Reichel]

DIAGONAL OF THE INVERSE

diag(inv(A)) in Dynamic Mean Field Theory (DMFT)

- Quantum mechanical studies of highly correlated particles
- Equation to be solved (repeatedly) is Dyson's equation

$$G(\omega) = [(\omega + \mu)I - V - \Sigma(\omega) + T]^{-1}$$

- ω (frequency) and μ (chemical potential) are real
 - V = trap potential = real diagonal
 - $\Sigma(\omega)$ == local self-energy - a complex diagonal
 - T is the hopping matrix (sparse real).
- Interested only in diagonal of $G(\omega)$ – in addition, equation must be solved self-consistently and ...
 - ... must do this for many ω 's

Stochastic Estimator

Notation:

- A = original matrix, $B = A^{-1}$.
- $\delta(B) = \text{diag}(B)$ [matlab notation]
- $\mathcal{D}(B)$ = diagonal matrix with diagonal $\delta(B)$
- \odot and \oslash : Elementwise multiplication and division of vectors
- $\{v_j\}$: Sequence of s random vectors

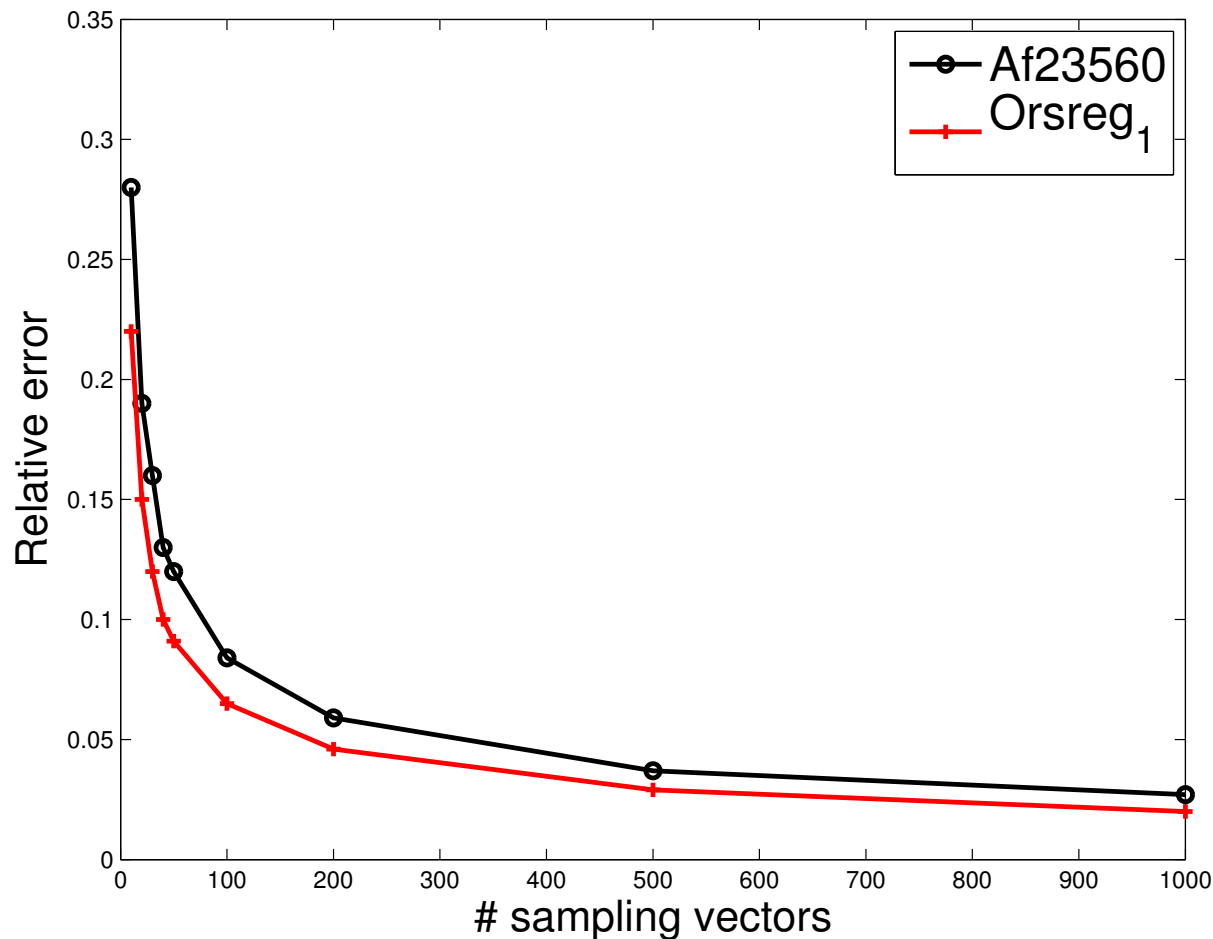
Result:

$$\delta(B) \approx \left[\sum_{j=1}^s v_j \odot B v_j \right] \oslash \left[\sum_{j=1}^s v_j \odot v_j \right]$$

Refs: C. Bekas , E. Kokiopoulou & YS ('05); C. Bekas, A. Curioni, I. Fedulova '09; ...

Typical convergence curve for stochastic estimator

- Estimating the diagonal of inverse of two sample matrices



➤ Let $V_s = [v_1, v_2, \dots, v_s]$. Then, alternative expression:

$$\mathcal{D}(B) \approx \mathcal{D}(BV_s V_s^\top) \mathcal{D}^{-1}(V_s V_s^\top)$$

Question: When is this result exact?

Answer:

- Let $V_s \in \mathbb{R}^{n \times s}$ with rows $\{v_{j,:}\}$; and $B \in \mathbb{C}^{n \times n}$ with elements $\{b_{jk}\}$
- Assume that: $\langle v_{j,:}, v_{k,:} \rangle = 0, \forall j \neq k, \text{ s.t. } b_{jk} \neq 0$

Then:

$$\mathcal{D}(B) = \mathcal{D}(BV_s V_s^\top) \mathcal{D}^{-1}(V_s V_s^\top)$$

➤ Approximation to b_{ij} exact when **rows** i and j of V_s are \perp

Using a sparse V : Probing

Goal:

Find V_s such that (1) s is small and (2) V_s satisfies Proposition (rows i & j orthogonal for any nonzero b_{ij})

Difficulty:

Can work only for sparse matrices but $B = A^{-1}$ is usually dense

- B can sometimes be approximated by a sparse matrix.
- Consider for some ϵ :
$$(B_\epsilon)_{ij} = \begin{cases} b_{ij}, & |b_{ij}| > \epsilon \\ 0, & |b_{ij}| \leq \epsilon \end{cases}$$
- B_ϵ will be sparse under certain conditions, e.g., when A is diagonally dominant
- In what follows we assume B_ϵ is sparse and set $B := B_\epsilon$.
- Pattern will be required by standard probing methods.

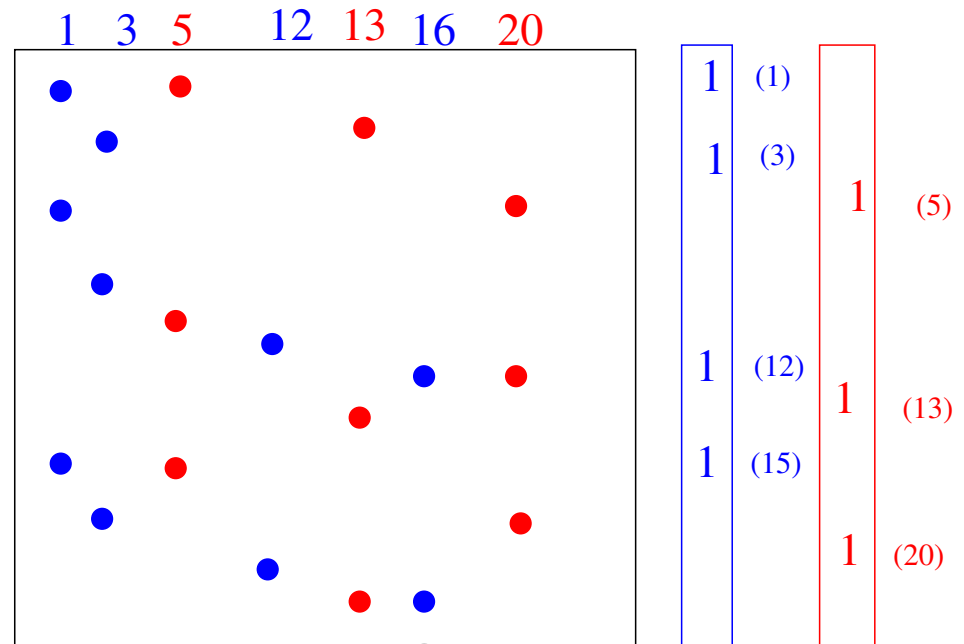
Standard probing (e.g. to compute a Jacobian)

- Several names for same method: “probing”; “CPR”, “Sparse Jacobian estimators”,...

Basis of the method: can compute Jacobian if a coloring of the columns is known so that no two columns of the same color overlap.

All entries of same color can be computed with one **matvec**.

Example: For all blue entries multiply B by the blue vector on right.



What about $\text{Diag}(\text{inv}(A))$?

- Define v_i - probing vector associated with color i :

$$[v_i]_k = \begin{cases} 1 & \text{if } \text{color}(k) == i \\ 0 & \text{otherwise} \end{cases}$$

- Standard probing satisfies requirement of Proposition but...
- ... this coloring is **not** what is needed! [It is an overkill]

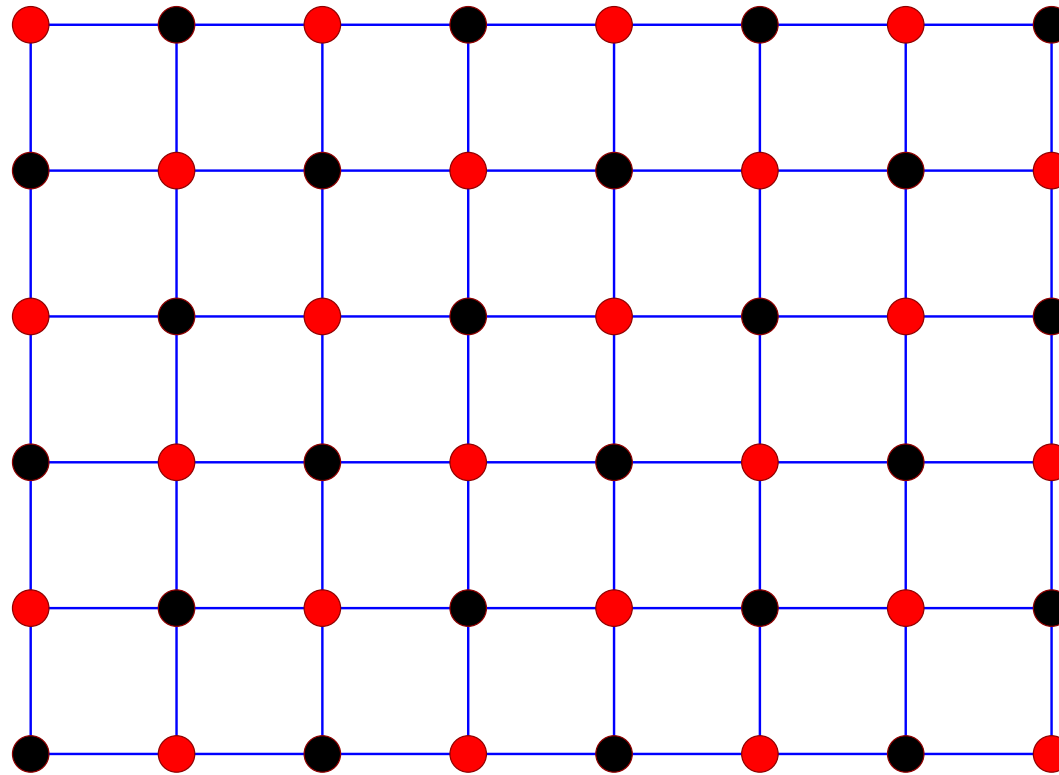
Alternative:

- Color the graph of B in the standard graph coloring algorithm [Adjacency graph, not graph of column-overlaps]

Result:

Graph coloring yields a valid set of probing vectors for $\mathcal{D}(B)$.

Example:



- Two colors required for this graph → two probing vectors
- Standard method: 6 colors [graph of $B^T B$]

Next Issue: Guessing the pattern of B

- Recall that we are dealing with $B := B_\epsilon$ ['pruned' B]
- Assume A diagonally dominant
- Write $A = D - E$, with $D = \mathcal{D}(A)$. Then :

$$A = D(I - F) \quad \text{with} \quad F \equiv D^{-1}E \quad \rightarrow$$

$$A^{-1} \approx \underbrace{(I + F + F^2 + \dots + F^k)}_{B^{(k)}} D^{-1}$$

- When A is D.D. $\|F^k\|$ decreases rapidly.
- Can approximate pattern of B by that of $B^{(k)}$ for some k .
- Interpretation in terms of paths of length k in graph of A .

Improvements

- Recent work by A. Stathopoulos, J. Laeuchli, and K. Orginos, on **hierarchical probing**. Produce approximate k -distance coloring of the graph to determine the patterns
- Somewhat specific to Lattice QCD
- E. Aune, D. P. Simpson, J. Eidsvik [Statistics and Computing 2012] combine probing with stochastic estimation. Good improvements reported.

EIGENVALUE COUNTS

Eigenvalue counts [with E. Polizzi and E. Di Napoli]

The problem:

- Find an **estimate** of the number of eigenvalues of a matrix in a given interval $[a, b]$.

Main motivation:

- Eigensolvers based on splitting the spectrum intervals and extracting eigenpairs from each interval independently.
- Contour integration-type methods:
 - FEAST approach [Polizzi 2011]
 - Sakurai-Sigiura method [2002]
- Polynomial filtering:
 - Schofield, Chelikowsky, YS'2011.

Eigenvalue counts: Standard approach

- Let spectrum of a Hermitian matrix A be

$$\lambda_1 \leq \lambda_2 \leq \cdots \leq \lambda_n$$

with eigenvectors u_1, u_2, \dots, u_n

- a, b such that $\lambda_1 \leq a \leq b \leq \lambda_n$.
- Want number $\mu_{[a,b]}$ of λ_i 's $\in [a, b]$
- Standard method: Use Sylvester inertia theorem
- Requires two LDL^T factorizations \rightarrow can be expensive!

- Alternative: Exploit trace of the eigen-projector:

$$P = \sum_{\lambda_i \in [a, b]} u_i u_i^T.$$

- We know that the trace of P is the wanted number $\mu_{[a,b]}$
- Goal: calculate an approximation to :

$$\mu_{[a,b]} = \text{Tr}(P).$$

- P is not available ... but can be approximated by
 - a polynomial in A , or
 - a rational function in A .

Eigenvalue counts: Approximation theory viewpoint

- Interpret P as a step function of A , namely:

$$P = h(A) \quad \text{where} \quad h(t) = \begin{cases} 1 & \text{if } t \in [a \ b] \\ 0 & \text{otherwise} \end{cases}$$

- Hutchinson's unbiased estimator uses only matrix-vector products to approximate the trace of a generic matrix A .
- Generate random vectors $v_k, k = 1, \dots, n_v$ with equally probable entries ± 1 . Then:

$$\text{tr}(A) \approx \frac{n}{n_v} \sum_{k=1}^{n_v} v_k^\top A v_k.$$

- No need to restrict values to ± 1

Polynomial filtering

- $h(t) \approx \psi(t)$, where ψ is a polynomial of degree k .
- We can estimate the trace of P as:

$$\mu_{[a,b]} \approx \frac{n}{n_v} \sum_{k=1}^{n_v} \mathbf{v}_k^\top \psi(\mathbf{A}) \mathbf{v}_k$$

- We use degree p Chebyshev polynomials:

$$h(t) \approx \psi_p(t) = \sum_{j=0}^p \gamma_j T_j(t).$$

Computing the polynomials: Jackson-Chebyshev

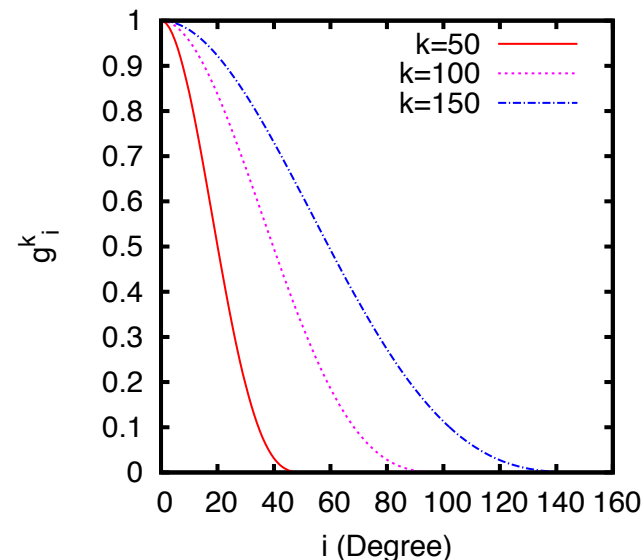
Chebyshev-Jackson approximation of a function f :

$$f(x) \approx \sum_{i=0}^k g_i^k \gamma_i T_i(x)$$

$$\gamma_i = \frac{2 - \delta_{i0}}{\pi} \int_{-1}^1 \frac{T_i(s)}{\sqrt{1-s^2}} f(s) ds \quad \delta_{i0} = \text{Kronecker symbol}$$

The g_i^k 's attenuate higher order terms in the sum.

Attenuation coefficient g_i^k for $k=50, 100, 150$ →



Let $\alpha_k = \frac{\pi}{k+2}$, then :

$$g_i^k = \frac{\left(1 - \frac{i}{k+2}\right) \sin(\alpha_k) \cos(i\alpha_k) + \frac{1}{k+2} \cos(\alpha_k) \sin(i\alpha_k)}{\sin(\alpha_k)}$$

See

Electronic structure calculations in plane-wave codes without diagonalization. Laurent O. Jay, Hanchul Kim, YS, and James R. Chelikowsky. *Computer Physics Communications*, 118:21–30, 1999.

The expansion coefficients γ_i

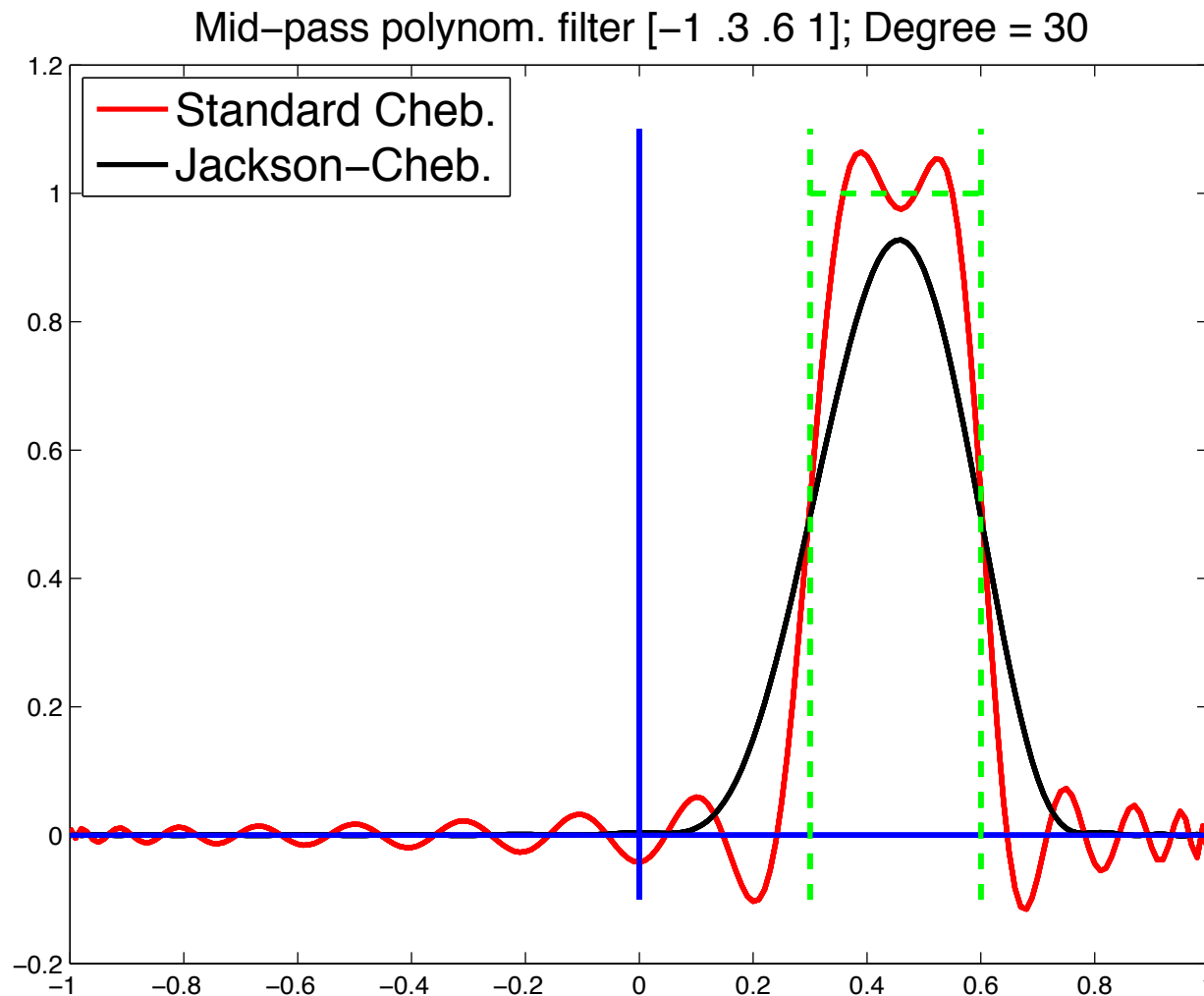
When $f(x)$ is a step function on $[a, b] \subseteq [-1, 1]$:

$$\gamma_i = \begin{cases} \frac{1}{\pi} (\arccos(a) - \arccos(b)) & : i = 0 \\ \frac{2}{\pi} \left(\frac{\sin(i \arccos(a)) - \sin(i \arccos(b))}{i} \right) & : i > 0 \end{cases}$$

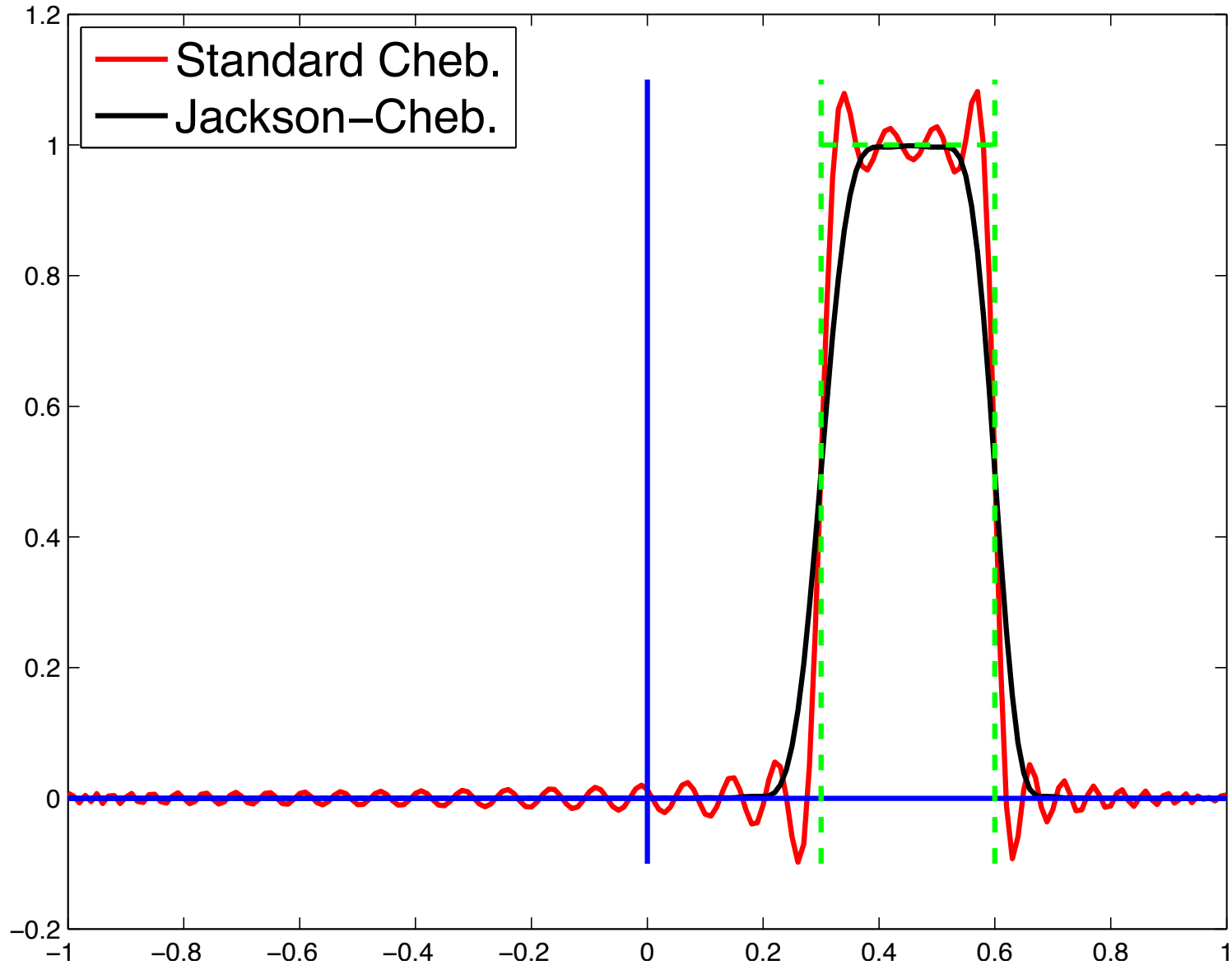
➤ A few examples follow –

Computing the polynomials: Jackson-Chebyshev

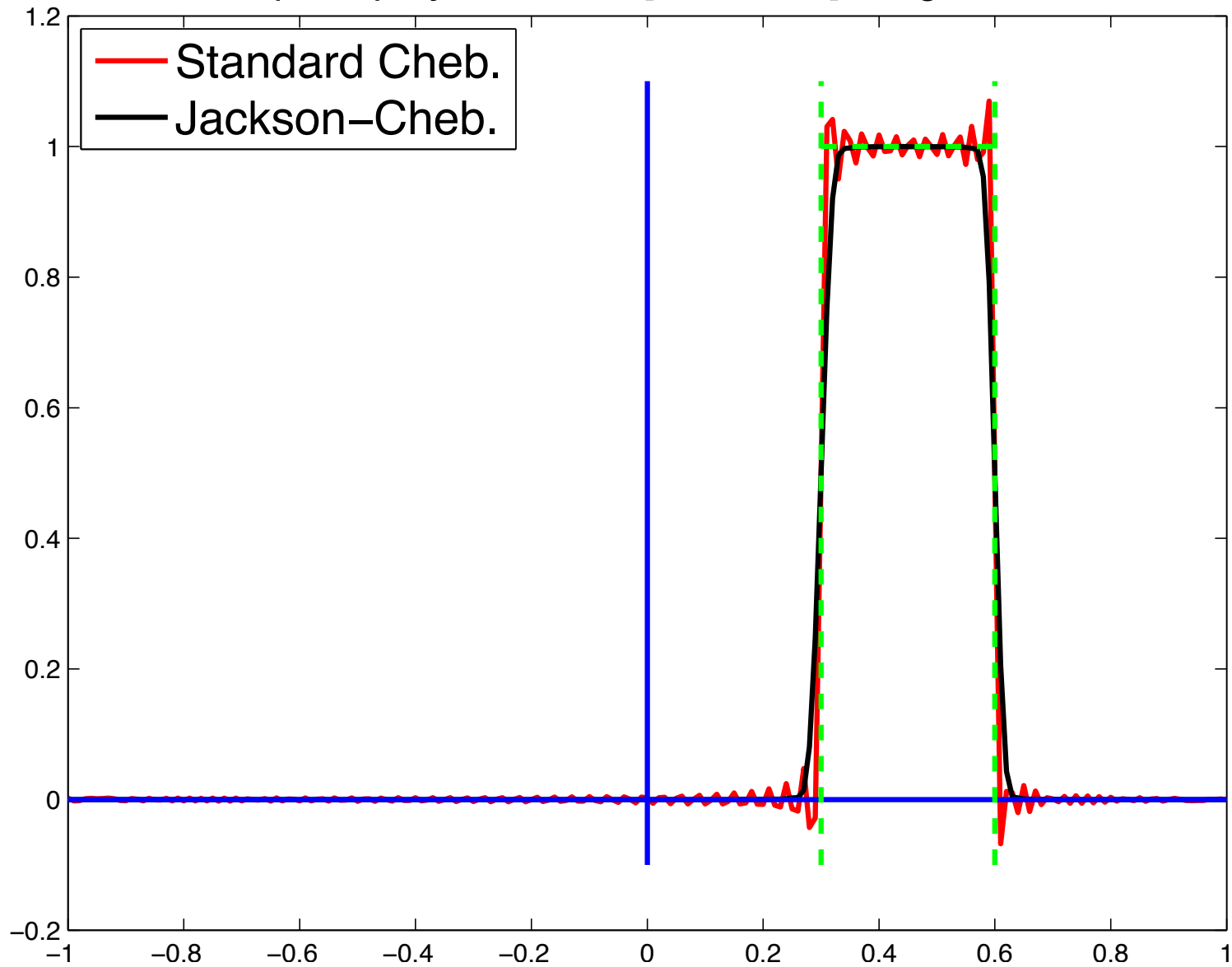
- Polynomials of degree 30 for $[a, b] = [.3, .6]$



Mid-pass polynom. filter [-1 .3 .6 1]; Degree = 80



Mid-pass polynom. filter [-1 .3 .6 1]; Degree = 200



$$\mu_{[a,b]} = \text{Tr}(P) \approx \frac{n}{n_v} \sum_{k=1}^{n_v} \left[\sum_{j=0}^p \gamma_j v_k^T T_j(A) v_k \right].$$

Easy to compute $T_j(A)v_k$ with 3-term recurrence of Chebyshev polynomials

$$w_{j+1} = 2Aw_j - w_{j-1}.$$

(A is transformed so its eigenvalues are in $[-1 \ 1]$)

Generalized eigenvalue problems

$$Ax = \lambda Bx$$

- Matrices A and B are symmetric and B is positive definite.

The projector P becomes

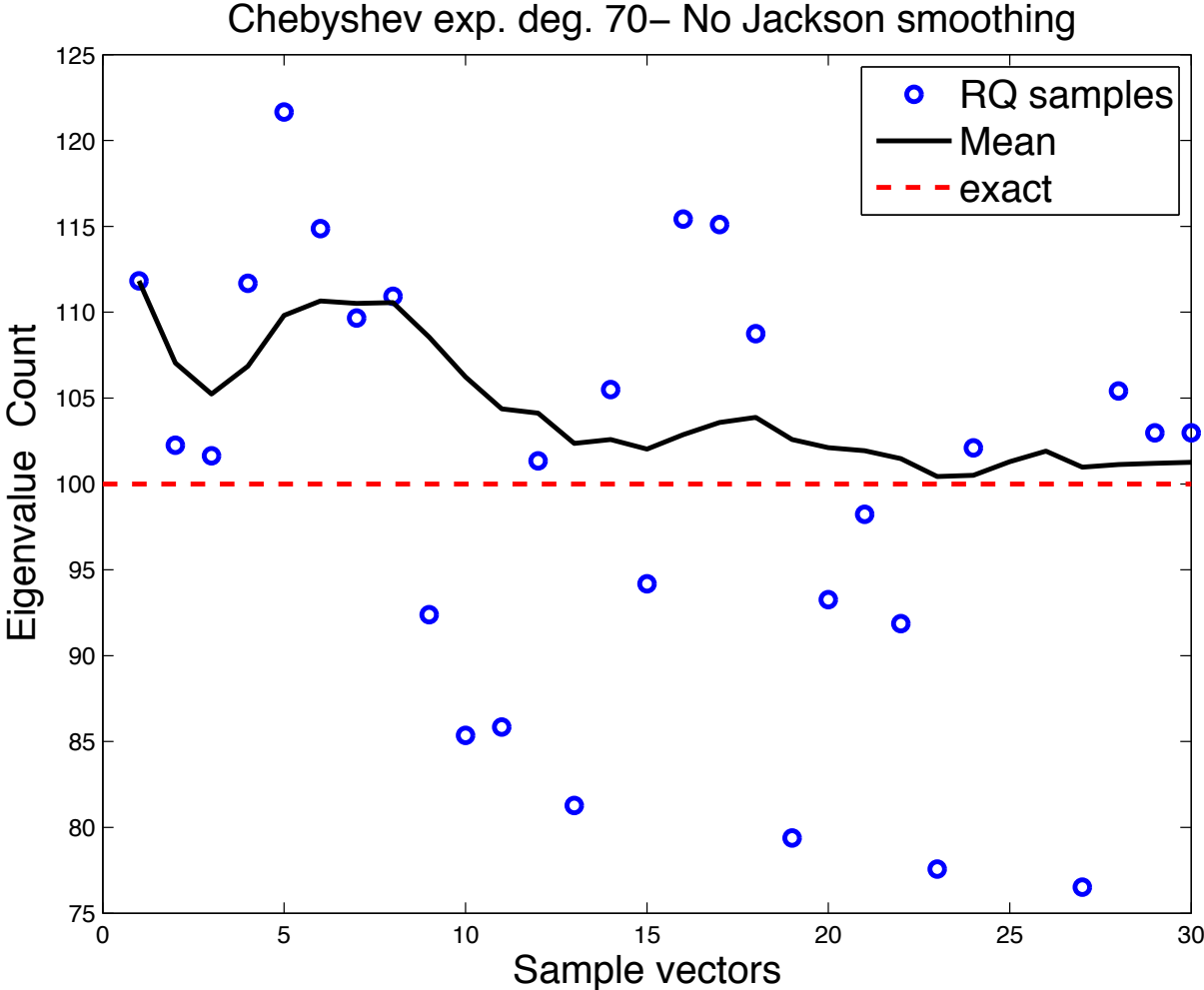
$$P = \sum_{\lambda_i \in [a \ b]} u_i u_i^T B,$$

- Again: Eigenvalue count == $\text{Tr}(P)$
- Exploit relation: $\text{inertia}(A - \sigma B) = \text{inertia}(B^{-1}A - \sigma I)$
- No need to factor or to solve systems

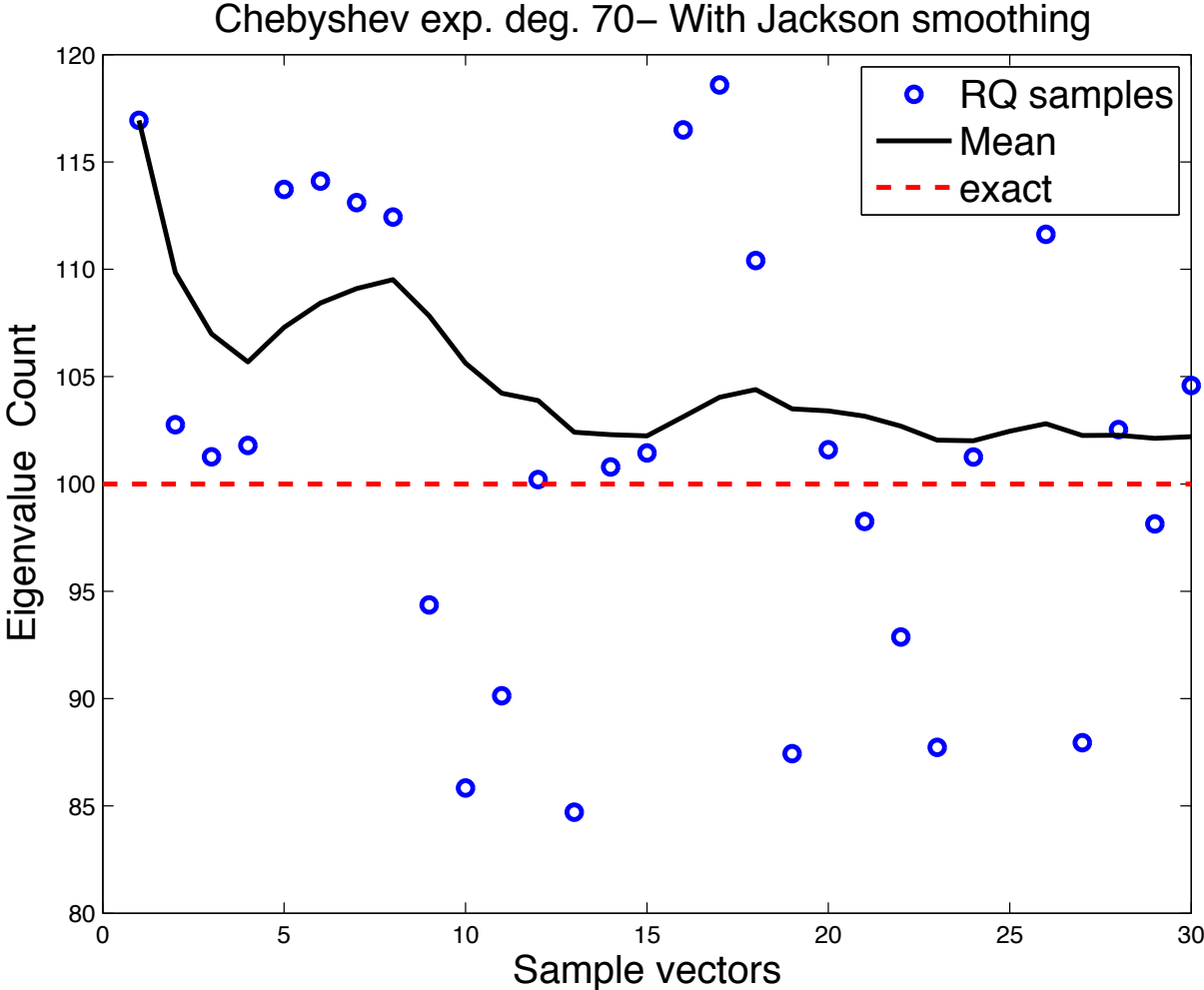
An example

- Matrix 'Na5' from PARSEC [see U. Florida collection]
- $n = 5832$, $nnz = 305630$ nonzero entries.
- Obtain the eigenvalue count when $a = (\lambda_{100} + \lambda_{101})/2$ and $b = (\lambda_{200} + \lambda_{201})/2$ so $\mu_{[a,b]} = 100$.
- Use pol. of degree 70.

Without Jackson Damping



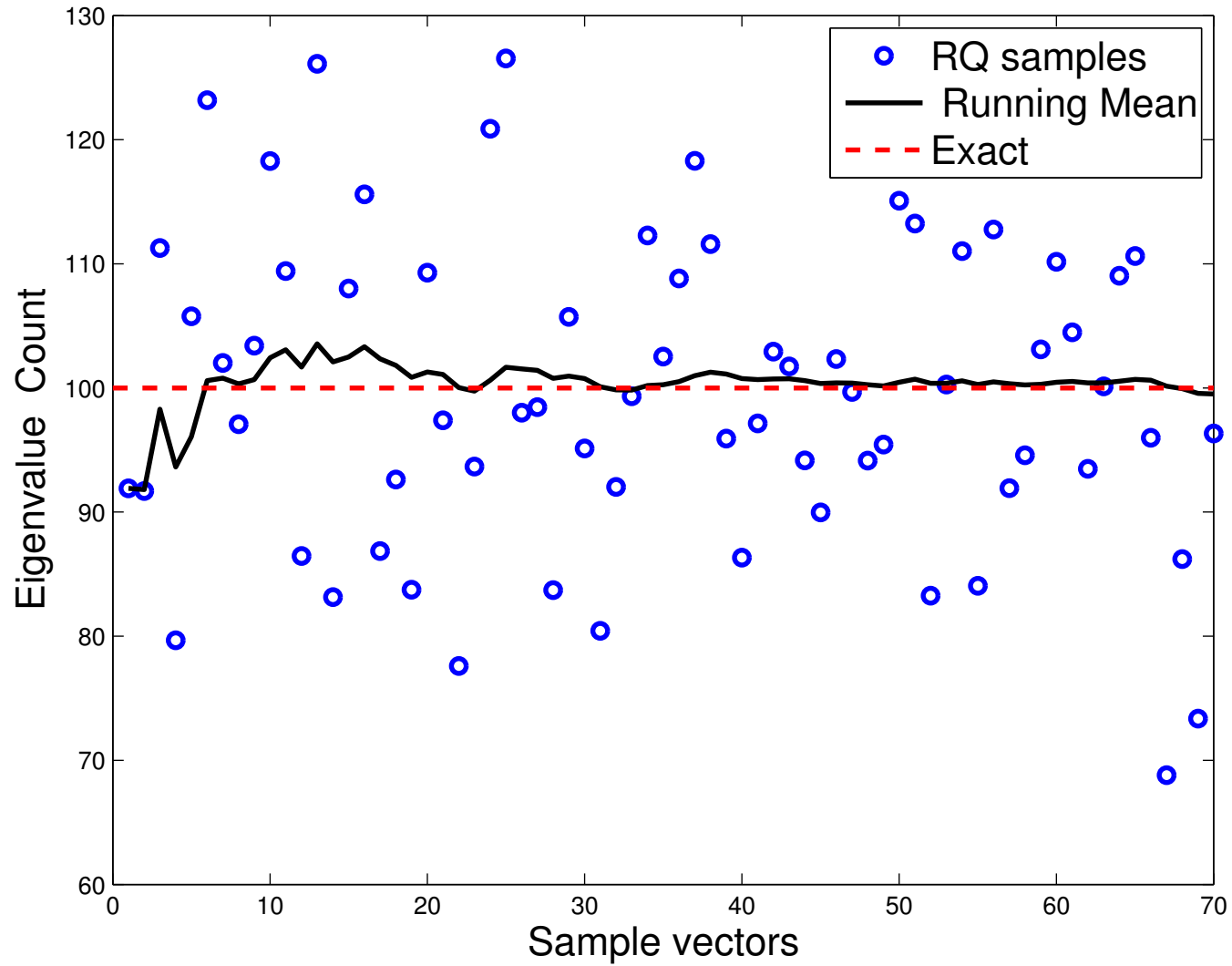
With Jackson Damping



An example from FEAST

- FEAST developed by Eric Polizzi (Amherst)..
- Based on a form of subspace iteration with a rational function of A
- Also works for generalized problems $Au = \lambda B$.
- Example: a small generalized problem ($n = 12,450$, $nnz = 86,808$).
- Result with standard Chebyshev shown. Deg=100, $nv = 70$.

Case: Gen2D; deg = 100; $n_v = 70$



➤ A few more comments:

- Method also works with rational approximations ...
- .. and it works for nonsymmetric problems (eigenvalues inside a given contour).
- For details see paper:

E. Di Napoli, E. Polizzi, and YS. *Efficient estimation of eigenvalue counts in an interval*. Preprint – see arXiv: <http://arxiv.org/abs/13>

DENSITY OF STATES

Computing Densities of States [with Lin-Lin and Chao Yang]

- Formally, the Density Of States (DOS) of a matrix A is

$$\phi(t) = \frac{1}{n} \sum_{j=1}^n \delta(t - \lambda_j),$$

where

- δ is the Dirac δ -function or Dirac distribution
 - $\lambda_1 \leq \lambda_2 \leq \dots \leq \lambda_n$ are the eigenvalues of A
- Note: number of eigenvalues in an interval $[a, b]$ is

$$\mu_{[a,b]} = \int_a^b \sum_j \delta(t - \lambda_j) dt \equiv \int_a^b n\phi(t) dt .$$

- $\phi(t)$ == a probability distribution function == probability of finding eigenvalues of A in a given infinitesimal interval near t .
- DOS is also referred to as the **spectral density**
- In Solid-State physics, λ_i 's represent single-particle energy levels.
- So the DOS represents # of levels per unit energy.
- Many uses in physics

Issue: How to deal with Distributions

- Highly ‘discontinuous’, not easy to handle numerically
- Solution for practical and theoretical purposes: replace ϕ by a ‘blurred’ (continuous) version ϕ_σ :

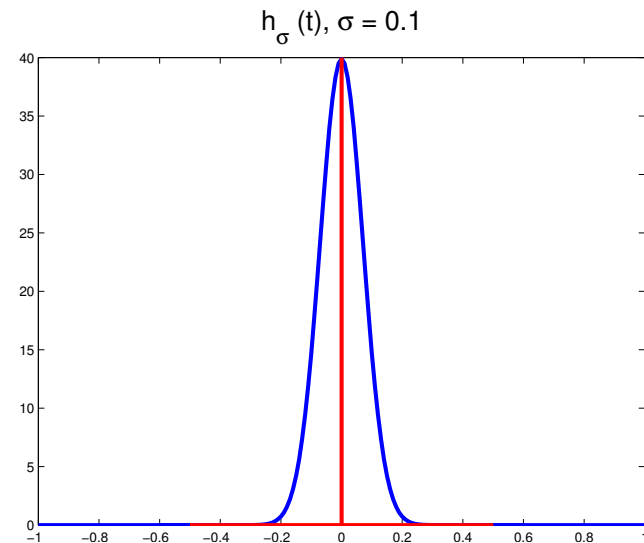
$$\phi_\sigma(t) = \frac{1}{n} \sum_{j=1}^n h_\sigma(t - \lambda_j),$$

where $h_\sigma(t) =$ any \mathcal{C}^∞ function s.t.:

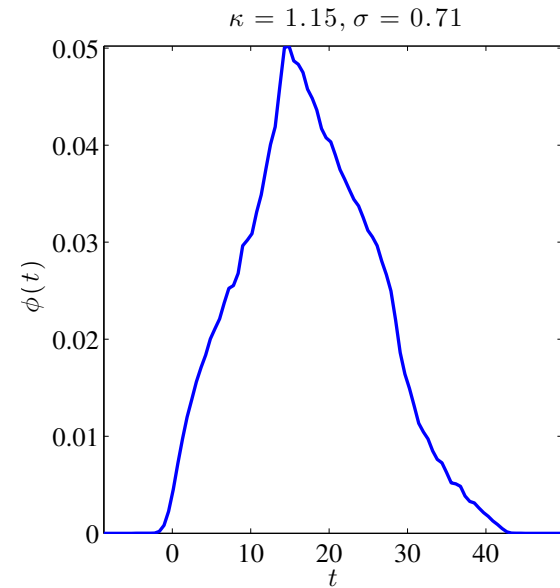
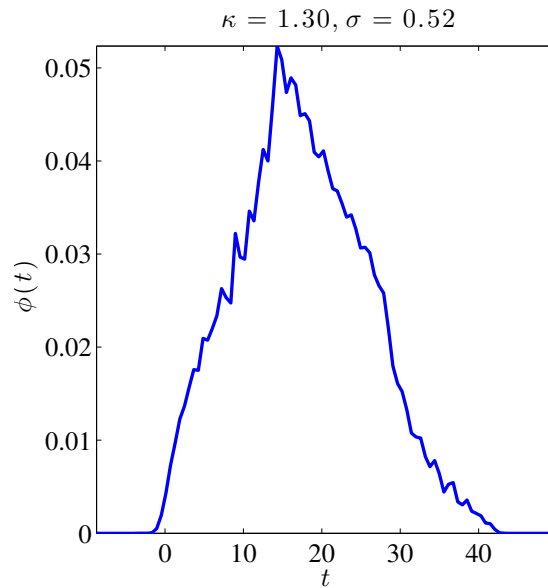
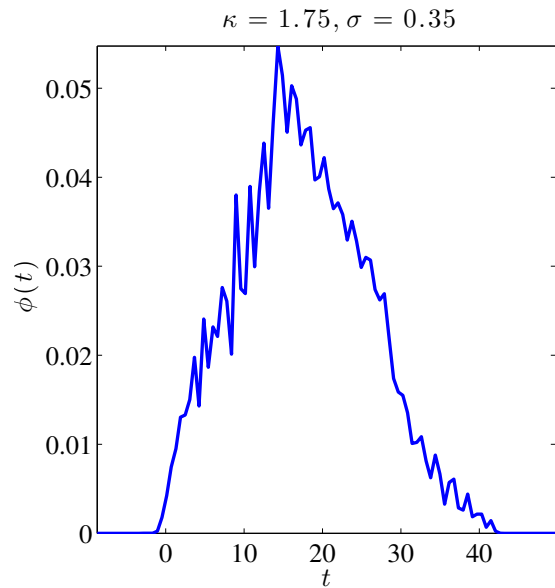
- $\int_{-\infty}^{+\infty} h_\sigma(s) ds = 1$
- h_σ has a peak at zero

- An example is the Gaussian:

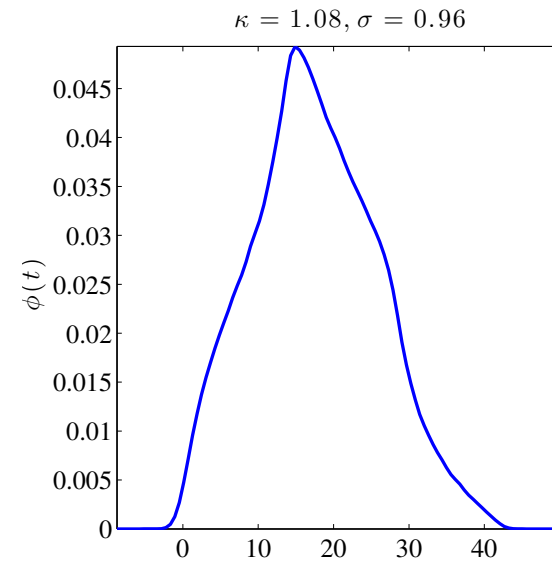
$$h_\sigma(t) = \frac{1}{(2\pi\sigma^2)^{1/2}} e^{-\frac{t^2}{2\sigma^2}}.$$



➤ How to select σ ? Example for Si_2



- Higher $\sigma \rightarrow$ smoother curve
- But loss of detail ..
- Compromise: $\sigma = \frac{h}{2\sqrt{2\log(\kappa)}}$,
- $h =$ resolution, $\kappa =$ parameter > 1



The Kernel Polynomial Method

- Used by Chemists to calculate the DOS – see Silver and Röder'94 , Wang '94, Drabold-Sankey'93, + others
- Basic idea: expand DOS into Chebyshev polynomials
- Use trace estimators [discovered independently] to get traces needed in calculations
- Assume change of variable done so eigenvalues lie in $[-1, 1]$.
- Include the weight function in the expansion so expand:

$$\hat{\phi}(t) = \sqrt{1-t^2}\phi(t) = \sqrt{1-t^2} \times \frac{1}{n} \sum_{j=1}^n \delta(t - \lambda_j).$$

Then, (full) expansion is: $\hat{\phi}(t) = \sum_{k=0}^{\infty} \mu_k T_k(t)$.

- Expansion coefficients μ_k are formally defined by:

$$\begin{aligned}\mu_k &= \frac{2 - \delta_{k0}}{\pi} \int_{-1}^1 \frac{1}{\sqrt{1-t^2}} T_k(t) \hat{\phi}(t) dt \\ &= \frac{2 - \delta_{k0}}{\pi} \int_{-1}^1 \frac{1}{\sqrt{1-t^2}} T_k(t) \sqrt{1-t^2} \phi(t) dt \\ &= \frac{2 - \delta_{k0}}{n\pi} \sum_{j=1}^n T_k(\lambda_j).\end{aligned}$$

- Here $2 - \delta_{k0} == 1$ when $k = 0$ and $== 2$ otherwise.
- Note: $\sum T_k(\lambda_i) = \text{Trace}[T_k(A)]$
- Estimate this, e.g., via stochastic estimator
- Generate random vectors $\mathbf{v}^{(1)}, \mathbf{v}^{(2)}, \dots, \mathbf{v}^{(n_{\text{vec}})}$
- Assume normal distribution with zero mean

- Each vector is normalized so that $\|v^{(l)}\| = 1, l = 1, \dots, n_{\text{vec}}$.
- Estimate the trace of $T_k(A)$ with stochastic estimator:

$$\text{Trace}(T_k(A)) \approx \frac{1}{n_{\text{vec}}} \sum_{l=1}^{n_{\text{vec}}} \left(v^{(l)}\right)^T T_k(A)v^{(l)}.$$

- Will lead to the desired estimate:

$$\mu_k \approx \frac{2 - \delta_{k0}}{n\pi n_{\text{vec}}} \sum_{l=1}^{n_{\text{vec}}} \left(v^{(l)}\right)^T T_k(A)v^{(l)}.$$

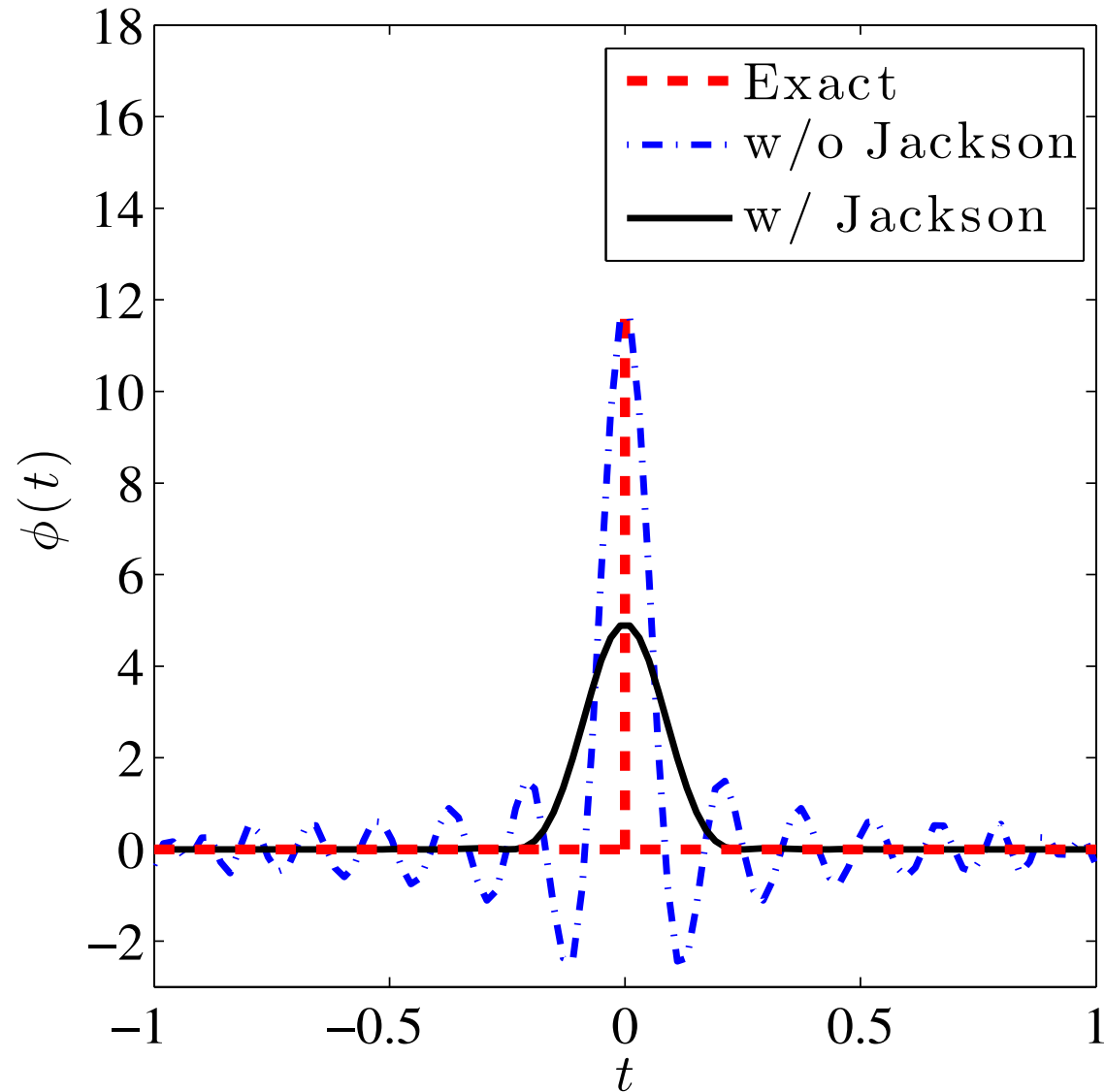
- To compute scalars of the form $v^T T_k(A)v$, exploit 3-term recurrence of the Chebyshev polynomial:

$$T_{k+1}(A)v = 2AT_k(A)v - T_{k-1}(A)v$$

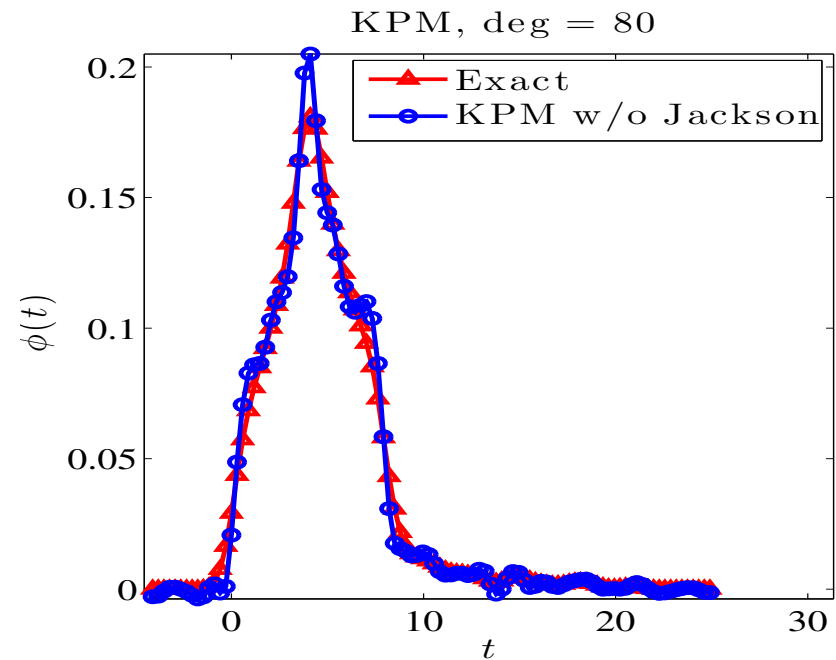
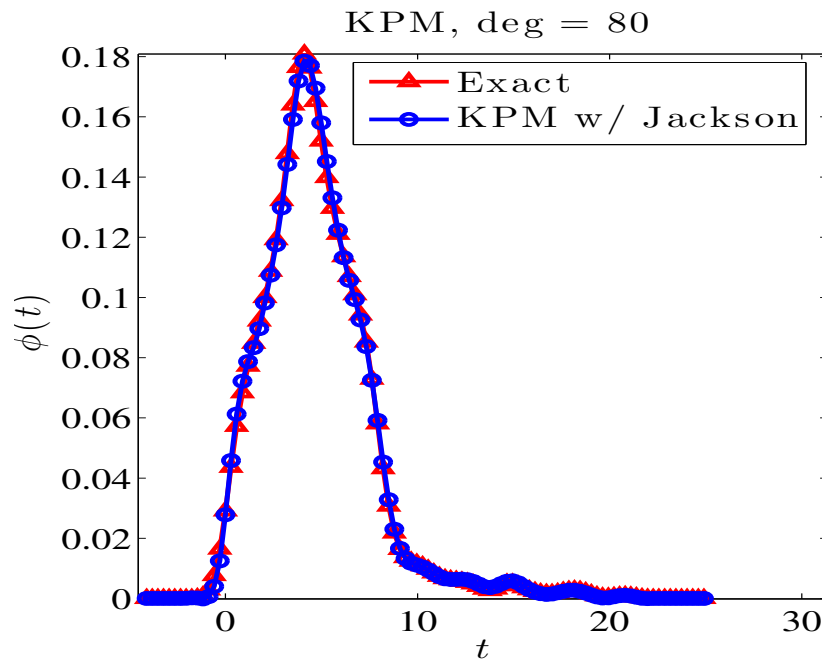
so if we let $v_k \equiv T_k(A)v$, we have

$$v_{k+1} = 2Av_k - v_{k-1}$$

- Same Jackson smoothing as before can be used



An example with degree 80 polynomials



Left: Jackson damping; right: without Jackson damping.

The Lanczos Spectroscopic approach

- Described in Lanczos' book "Applied Analysis, (1956)" as a means to compute eigenvalues.
- Idea: assimilate λ_i 's to frequencies and perform Fourier analysis to extract them
- Also relies on Chebyshev polynomials
- Though not emphasized in the description, the method uses random sampling
- Let B a symmetric real matrix with eigenvalues in $[-1,1]$
- Let v_0 == an initial vector – expand in eigenbasis as

$$v_0 = \sum_{j=1}^n \beta_j u_j, \quad \text{with} \quad \beta_j = u_j^T v_0$$

- Let $v_k = T_k(A)v_0$, for $k = 0, \dots, M$. Then:

$$v_0^T v_k = \sum_{j=1}^n \beta_j^2 T_k(\lambda_j) = \sum_{j=1}^n \beta_j^2 \cos(k\theta_j), \text{ with } \lambda_j = \cos \theta_j.$$

View $v_0^T v_k$ as a discretization of the **periodic** function to the right sampled at $t = 0, 1, \dots, M$.

$$f(t) = \sum_{j=1}^n \beta_j^2 \cos(t\theta_j)$$

- Problem: find values of θ_j , for $j = 1, \dots, n$
- Compute cosine transform of f ; For $p = 0, \dots, M$:

$$F(p) = \frac{f(0) + (-1)^p f(M)}{2} + \sum_{k=1}^{M-1} f(k) \cos \frac{kp\pi}{M},$$

- If f has an eigenvalue $\lambda = \cos \theta$, then component $\cos(\theta t)$, revealed by a peak at the point

$$p = \frac{l\theta}{\pi}.$$

- Peak at p_j corresponds to eigenvalue $\lambda_j = \cos \theta_j$ with $\theta_j = (p_j/M)\pi$, and so,

$$\lambda_j = \cos(\theta_j) = \cos(p_j\pi/M)$$

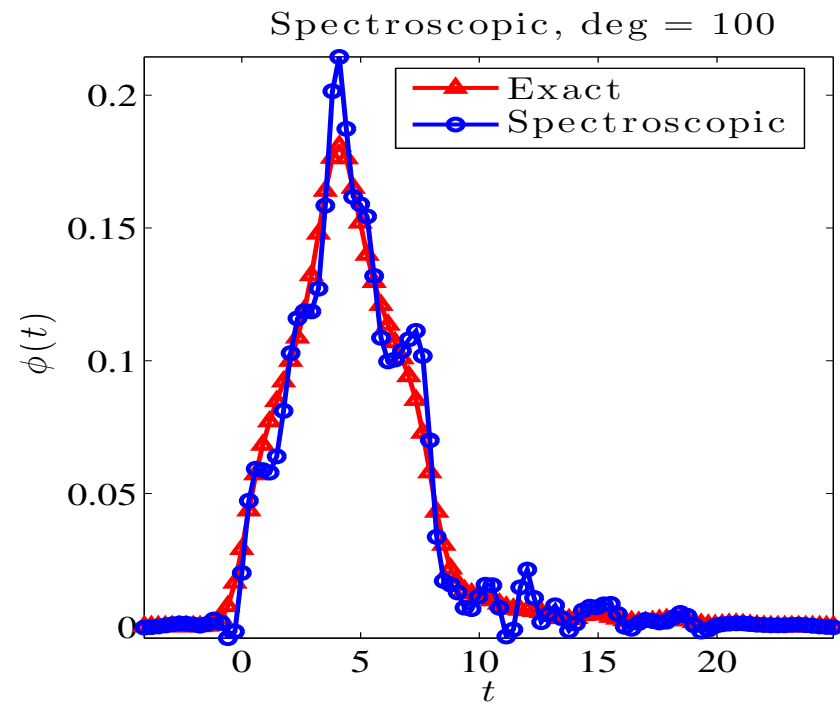
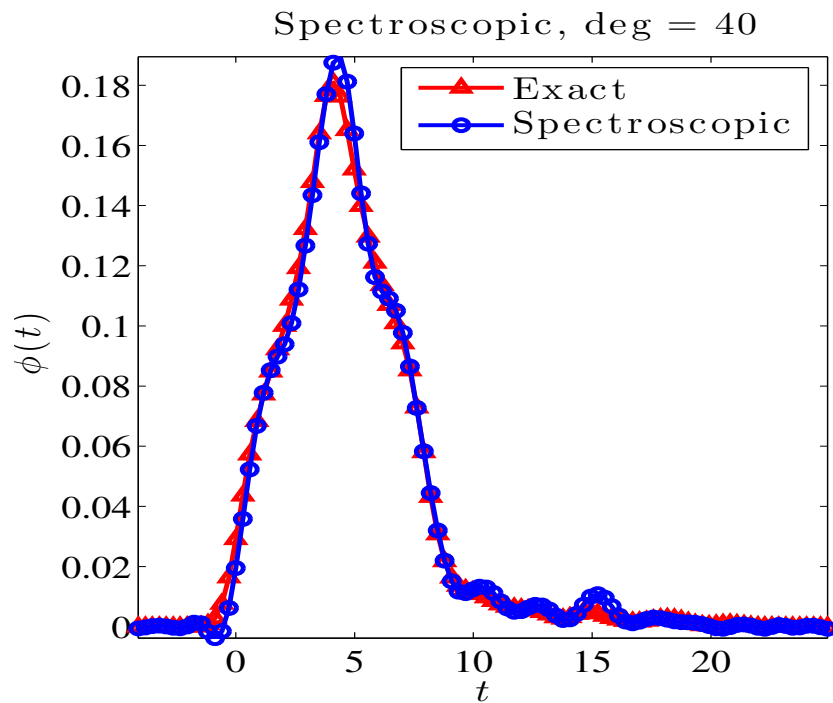
- For a sequence of random vectors compute

$$\hat{F}(\hat{p}) \equiv F\left(\frac{M}{\pi} \arccos \hat{p}\right), \quad \hat{p} = \cos(p\pi/M), p = 0 : M.$$

- Average these values $\rightarrow \phi(t_i) \approx Cst \times \hat{F}(t_i)$

The Lanczos Spectroscopic approach: Example

- Same example as before



Left: Degree 40; Right: degree 100

Recall: How to deal with Distributions

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- Solution for practical and theoretical purposes: replace ϕ by a ‘blurred’ (continuous) version ϕ_σ :

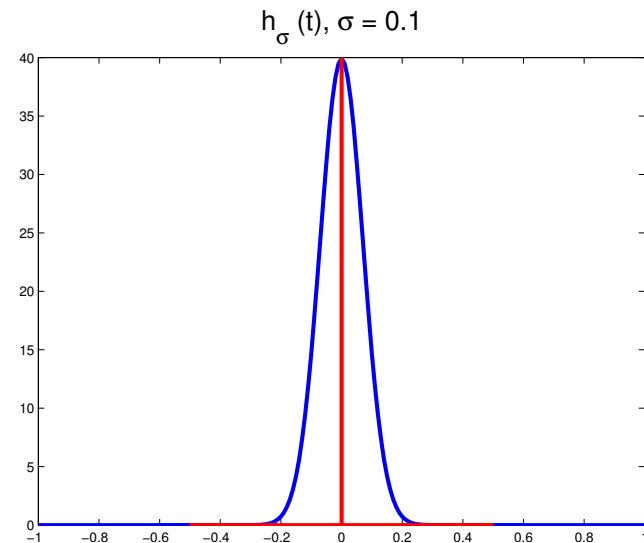
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- h_σ has a peak at zero

- An example is the Gaussian:

$$h_\sigma(t) = \frac{1}{(2\pi\sigma^2)^{1/2}} e^{-\frac{t^2}{2\sigma^2}}.$$



Delta-Gauss Legendre

- Idea: Instead of approximating ϕ directly, first select a representative ϕ_σ of ϕ for a given σ and then approximate ϕ_σ .
- ϕ_σ is a 'surrogate' for ϕ . Obtained by replacing δ_λ by :

$$h_\sigma(\lambda - t) = \frac{1}{(2\pi\sigma^2)^{1/2}} \exp \left[-\frac{(\lambda - t)^2}{2\sigma^2} \right].$$

- Goal: to expand into Legendre polynomials $L_k(\lambda)$
- With normalization factor expansion is written as:

$$h_\sigma(\lambda - t) = \frac{1}{(2\pi\sigma^2)^{1/2}} \sum_{k=0}^{\infty} \left(k + \frac{1}{2} \right) \gamma_k L_k(\lambda) .$$

- To determine the γ_k 's we will also need to compute:

$$\psi_k = \int_{-1}^1 L'_k(s) e^{-\frac{1}{2}((s-t)/\sigma)^2} ds.$$

Set $\zeta_k = e^{-\frac{1}{2}((1-t)/\sigma)^2} - (-1)^k e^{-\frac{1}{2}((1+t)/\sigma)^2}$.

- Then, for $k = 0, 1, \dots$,:

$$\begin{cases} \gamma_{k+1} = \frac{2k+1}{k+1} [\sigma^2(\psi_k - \zeta_k) + t\gamma_k] - \frac{k}{k+1}\gamma_{k-1} \\ \psi_{k+1} = (2k+1)\gamma_k + \psi_{k-1}. \end{cases}$$

Initialization: set $\gamma_{-1} = \psi_{-1} = 0$ $\psi_1 = \gamma_0$, and $\psi_0 = 0$ and:

$$\gamma_0 = \sigma \sqrt{\frac{\pi}{2}} \left[\operatorname{erf} \left(\frac{1-t}{\sqrt{2}\sigma} \right) + \operatorname{erf} \left(\frac{1+t}{\sqrt{2}\sigma} \right) \right],$$

Use of the Lanczos Algorithm

- Background: The Lanczos algorithm generates an orthonormal basis $V_m = [v_1, v_2, \dots, v_m]$ for the **Krylov subspace**:

$$\text{span}\{v_1, Av_1, \dots, A^{m-1}v_1\}$$

- ... such that:

$$V_m^H AV_m = T_m \text{ - with}$$

$$T_m = \begin{pmatrix} \alpha_1 & \beta_2 & & & & \\ & \beta_2 & \alpha_2 & \beta_3 & & \\ & & \beta_3 & \alpha_3 & \beta_4 & \\ & & & \cdot & \cdot & \cdot \\ & & & & \cdot & \cdot \\ & & & & & \beta_m & \alpha_m \end{pmatrix}$$

- Lanczos builds orthogonal polynomials wrt to dot product:

$$\int p(t)q(t)dt \equiv (p(A)v_1, q(A)v_1)$$

- In theory v_i 's defined by 3-term recurrence are orthogonal.
- Let $\theta_i, i = 1 \dots, m$ be the eigenvalues of T_m [Ritz values]
- y_i 's associated eigenvectors; Ritz vectors: $\{V_m y_i\}_{i=1:m}$
- Ritz values approximate eigenvalues [from 'outside in']
- Could compute θ_i 's then get approximate DOS from these
- Problem: θ_i not good enough approximations – especially inside the spectrum.

- Better idea: exploit relation of Lanczos with (discrete) orthogonal polynomials and related Gaussian quadrature:

$$\int p(t) dt \approx \sum_{i=1}^m a_i p(\theta_i) \quad a_i = [e_1^T y_i]^2$$

- See, e.g., Golub & Meurant '93, and also Gautschi'81, Golub and Welsch '69.
- Formula exact when p is a polynomial of degree $\leq 2m + 1$

- Consider now $\int p(t)dt =$ discrete (Stieljes) integral \equiv

$$(p(A)v, v) = \sum \beta_i^2 p(\lambda_i) \equiv \langle \phi_v, p \rangle$$

- Then $\langle \phi_v, p \rangle \approx \sum a_i p(\theta_i) = \sum a_i \langle \delta_{\theta_i}, p \rangle \rightarrow$

$$\phi_v \approx \sum a_i \delta_{\theta_i}$$

- To mimick the effect of $\beta_i = 1, \forall i$, use several vectors v and average the result of the above formula over them..

Experiments

- Goal: to compare errors for similar number of matrix-vector products
- Example: Kohn-Sham Hamiltonian associated with a benzene molecule generated from PARSEC; size $n = 8,219$
- In all cases, we use 10 sampling vectors
- General observation: DGL, Lanczos, and KPM are best,
- Spectroscopic method does OK
- Haydock's method [another method based on the Lanczos algorithm] not as good

Method	L^1 error	L^2 error	L^∞ error
KPM w/ Jackson, deg=80	2.592e-02	5.032e-03	2.785e-03
KPM w/o Jackson, deg=80	2.634e-02	4.454e-03	2.002e-03
KPM Legendre, deg=80	2.504e-02	3.788e-03	1.174e-03
Spectroscopic, deg=40	5.589e-02	8.652e-03	2.871e-03
Spectroscopic, deg=100	4.624e-02	7.582e-03	2.447e-03
DGL, deg=80	1.998e-02	3.379e-03	1.149e-03
Lanczos, deg=80	2.755e-02	4.178e-03	1.599e-03
Haydock, deg=40	6.951e-01	1.302e-01	6.176e-02
Haydock, deg=100	2.581e-01	4.653e-02	1.420e-02

L^1 , L^2 , and L^∞ error compared with the normalized “surrogate” DOS for benzene matrix

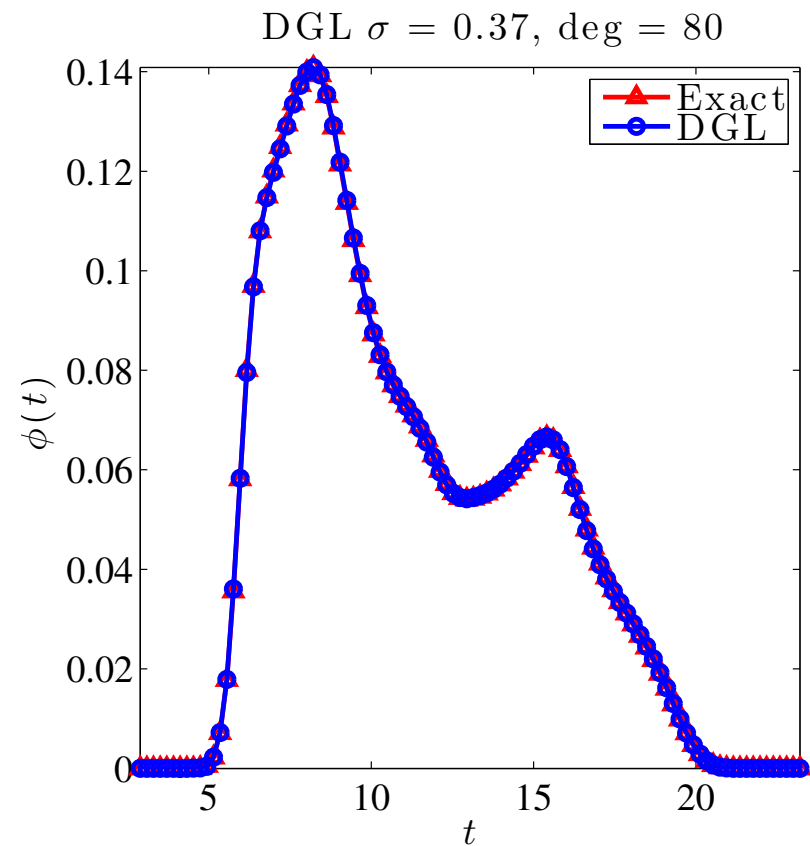
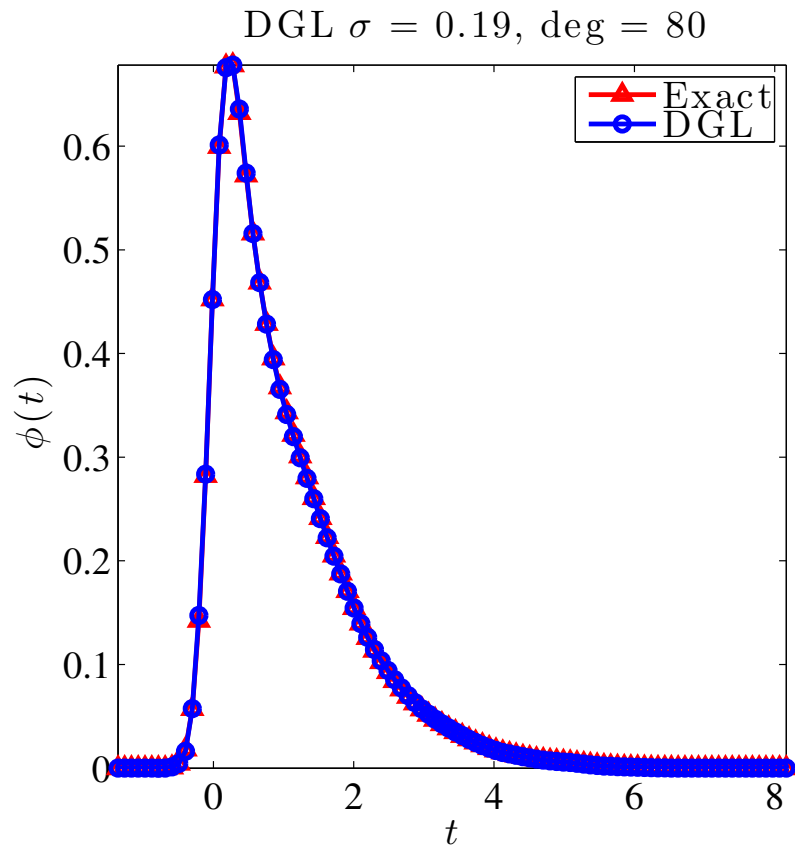
Other matrices

Matrix	n	λ_1	λ_n
Ga ₁₀ As ₁₀ H ₃₀	113,081	-1.2	1.3×10^3
PE3K	9,000	8.1×10^{-6}	1.3×10^2
CFD1	70,656	2.0×10^{-5}	6.8
SHWATER	81,920	5.8	2.0×10^1

Description of the size and the spectrum range of the test matrices.

Matrix	Method	L^1 error	L^2 error	L^∞ error
Ga ₁₀ As ₁₀ H ₃₀	DGL	3.937e-03	3.214e-04	4.301e-05
	Lanczos	4.828e-03	3.940e-04	5.452e-05
PE3K	DGL	4.562e-03	7.368e-04	3.143e-04
	Lanczos	5.459e-03	7.372e-04	3.294e-04
CFD1	DGL	2.276e-03	1.299e-03	1.746e-03
	Lanczos	2.024e-03	1.286e-03	2.478e-03
SHWATER	DGL	3.779e-03	1.282e-03	9.328e-04
	Lanczos	3.047e-03	9.829e-04	6.100e-04

L^1 , L^2 , and L^∞ error associated with the approximate spectral densities produced by the DGL and Lanczos methods for different test matrices.



Approximate spectral densities of CFD1 and SHWATER matrices obtained by DGL along with exact smoothed ones

Conclusion

- Probabilistic algorithms provide powerful tools for solving various problems: eigenvalue counts, DOS, $\text{Diag}(f(A))$..
- Most of the algorithms we discussed rely on estimating trace of $f(A)$.

Q: Can we do better than random sampling [e.g., probing,..]?

- Physicists are interested in modified forms of the density of states. → Explore extensions of what we did.